



Counterflow of impurities in harmonically confined optical lattices

Phys. Rev. Lett. 135, 023404 (2025)

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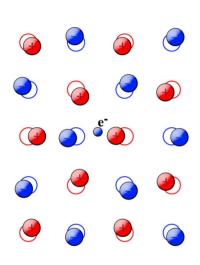
QUOST IX, Valparaiso, 7 November 2025

Collaborators

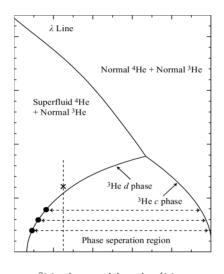
A. Rojo-Francàs, B. Juliá-Díaz Universitat de Barcelona L. Morales-Molina PUC Chile

Impurities

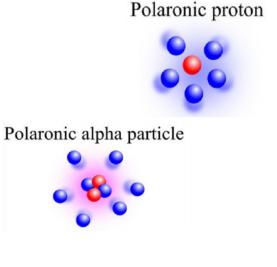
 The study of impurities immersed in a quantum medium has a long history and is relevant in many fields of physics.



Electrons in an ionic crystal L. Landau and S. Pekar, Zh. Eksp. Teor. Fiz **18**, 419 (1948).



³He impurities in ⁴He G. Baym and C. Pethick, "Landau Fermi-Liquid Theory: Concepts and Application" (1991).



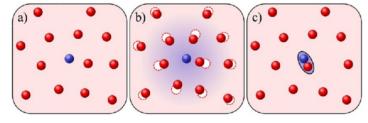
Impurities in nuclear systems Tajima *et al.*, AAPPS Bulletin, **34**, 9 (2024).

Impurities in ultracold atomic gases

 The study of impurities has been revitalised thanks to their experimental realisation in ultracold atomic mixtures with a high population imbalance.

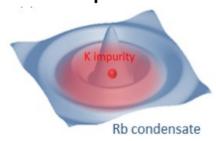
F. Grusdt et al. Rep. Prog. Phys. 88, 066401 (2025). P. Massignan et al., arXiv:2501.09618 (2025),

Fermi polaron



A. Schirotzek et al., PRL 102, 230402 (2009).

Bose polaron



M.-G. Hu et al., PRL 117, 055301 (2016).

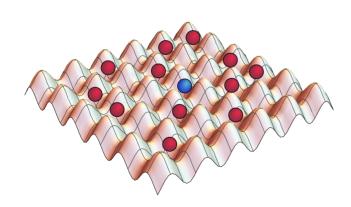
 Ultracold atoms offer a high level of controllability, making them an ideal platform for studying impurities.

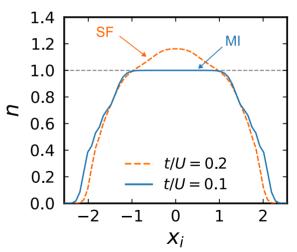
Impurities in optical lattices

- A rich platform for studying ultracold atomic impurities is tight optical lattices.

 M. Lewenstein et al., Ultracold Atoms in Optical Lattices: Simulating Quantum Many-Body Systems (Oxford, 2012).
- Bosons in optical lattices feature a superfluid-to-Mott insulator (SF-MI) transition.
- The study of impurities across this SF-MI transition has attracted increasing attention.

V. E. Colussi, F. Caleffi, C. Menotti, and A. Recati, Phys. Rev. Lett 130, 17, 3002 (2023).





Density profiles of bosons in a one-dimensional optical lattice with harmonic confinement.

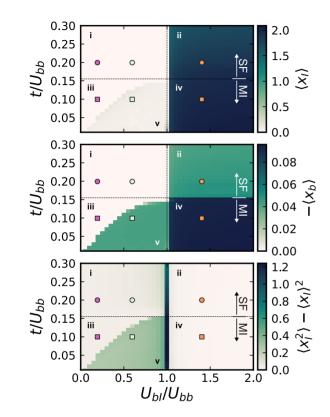
Impurities in harmonically confined optical lattices

 We studied a single impurity interacting with a bosonic bath in a onedimensional harmonically confined optical lattice.

$$\begin{split} \hat{H} &= -t \sum_{\sigma = b, I} \sum_{i} \left(\hat{a}_{i,\sigma}^{\dagger} \hat{a}_{i+1,\sigma} + \text{h.c.} \right) + \underbrace{V_{\text{ho}} \sum_{i,\sigma = b, I} i^2 \hat{n}_{i,\sigma}}_{\text{Harmonic trap}} \\ &+ \underbrace{\frac{U_{bb}}{2} \sum_{i} \hat{n}_{i,b} \left(\hat{n}_{i,b} - 1 \right)}_{i} + \underbrace{U_{bI} \sum_{i} \hat{n}_{i,b} \hat{n}_{i,I}}_{\text{Bath-impurity repulsion}} \\ &+ \underbrace{U_{bb} \sum_{i} \hat{n}_{i,b} \left(\hat{n}_{i,b} - 1 \right)}_{\text{Bath-impurity repulsion}} + \underbrace{U_{bb}, U_{bI} > 0}_{\text{Bath-impurity repulsion}} \end{split}$$

 We have investigated this system using the density matrix renormalisation group (DMRG) for large lattices with a large number of particles.

S. R. White, Phys. Rev. Lett. 69, 2863 (1992). U. Schollwöck, Rev. Mod. Phys. 77, 259 (2005).



Density profiles:

$$n_{\sigma}(i) = \langle \hat{n}_{i,\sigma} \rangle$$

Average position:

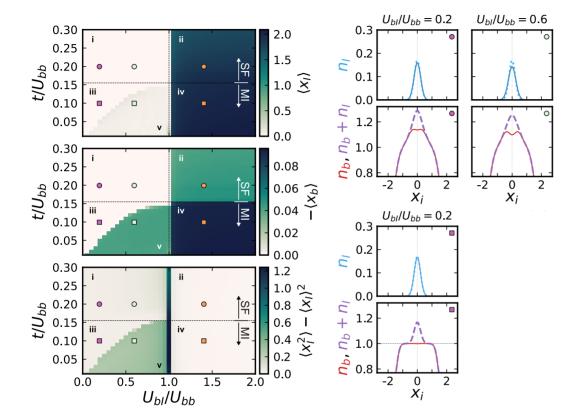
$$\langle x_{\sigma} \rangle = \frac{1}{N_{\sigma}} \sum_{i} x_{i} n_{\sigma}(i)$$

Impurity cloud:

$$\langle x_I^2 \rangle = \frac{1}{N_\sigma} \sum_i x_i^2 n_I(i)$$

$$x_i = i d/\xi, \quad \xi = d\sqrt{t/V_{\text{ho}}}.$$

The system shows well-defined phases.



Density profiles:

$$n_{\sigma}(i) = \langle \hat{n}_{i,\sigma} \rangle$$

Average position:

$$\langle x_{\sigma} \rangle = \frac{1}{N_{\sigma}} \sum_{i} x_{i} n_{\sigma}(i)$$

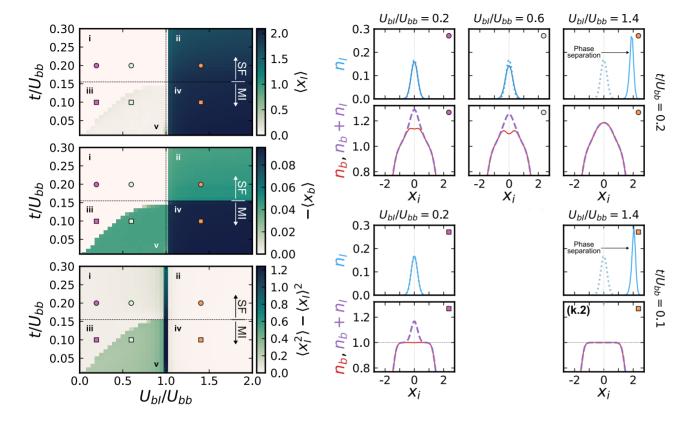
Impurity cloud:

 $t/U_{bb} = 0.1$

$$\langle x_I^2 \rangle = \frac{1}{N_\sigma} \sum_i x_i^2 n_I(i)$$

$$x_i = i d/\xi, \quad \xi = d\sqrt{t/V_{\text{ho}}}.$$

Regions i and iii are miscible phases.



Density profiles:

$$n_{\sigma}(i) = \langle \hat{n}_{i,\sigma} \rangle$$

Average position:

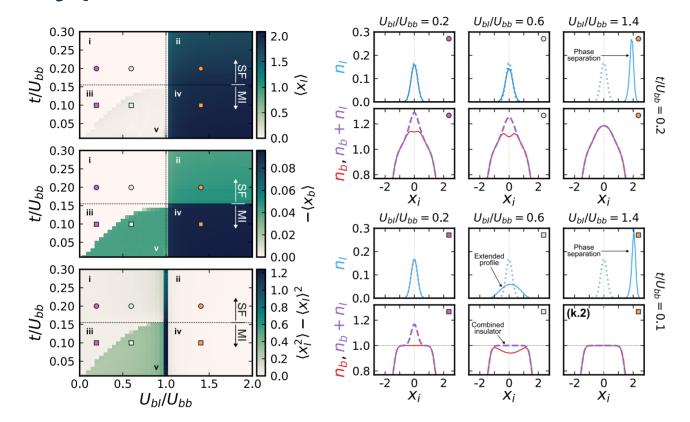
$$\langle x_{\sigma} \rangle = \frac{1}{N_{\sigma}} \sum_{i} x_{i} n_{\sigma}(i)$$

Impurity cloud:

$$\langle x_I^2 \rangle = \frac{1}{N_\sigma} \sum_i x_i^2 n_I(i)$$

$$x_i = i d/\xi, \quad \xi = d\sqrt{t/V_{\text{ho}}}.$$

Regions ii and iv are phase-separated configurations.



Density profiles:

$$n_{\sigma}(i) = \langle \hat{n}_{i,\sigma} \rangle$$

Average position:

$$\langle x_{\sigma} \rangle = \frac{1}{N_{\sigma}} \sum_{i} x_{i} n_{\sigma}(i)$$

Impurity cloud:

$$\langle x_I^2 \rangle = \frac{1}{N_\sigma} \sum_i x_i^2 n_I(i)$$

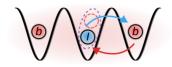
$$x_i = i d/\xi, \quad \xi = d\sqrt{t/V_{\text{ho}}}.$$

Region v is a phase where the impurity shows an extended profile, and the
combined profile of bath and impurity displays a domain of unity filling.

Counterflow

 The new state corresponds to a counterflow phase with long-range anti-pair order,

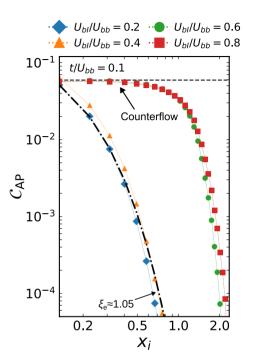
$$C_{\rm AP} = \langle \hat{a}_{0,I} \hat{a}_{0,b}^{\dagger} \hat{a}_{i,b}^{\dagger} \hat{a}_{i,I} \rangle.$$

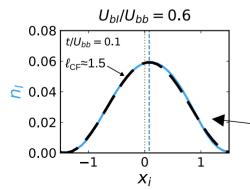


 Supercounterflows were realised experimentally very recently with binary Mott insulators.

Y.-G. Zheng et al., Nat. Phys. 21, 208 (2025).

 Our results show that counterflows appear for a large population imbalance.





- The counterflow features an orthogonality catastrophe.
- The impurity behaves as a free particle in a square well,

$$n_I^{(CF)}(x_i) = n_i^{(0)} \cos^2(\pi(x_i - \langle x_I \rangle)/\ell_{CF}).$$

Conclusions

- An impurity can form a non-trivial counterflow state with a bosonic bath in harmonically confined optical lattices.
- This means that counterflows form in mixtures with high population imbalances.
- Future work:
 - → Dynamics.
 - → Multiple impurities.
 - → Bose-Bose and Bose-Fermi mixtures.



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